



TRANSVERSE ENTANGLEMENT OF BIPHOTONS

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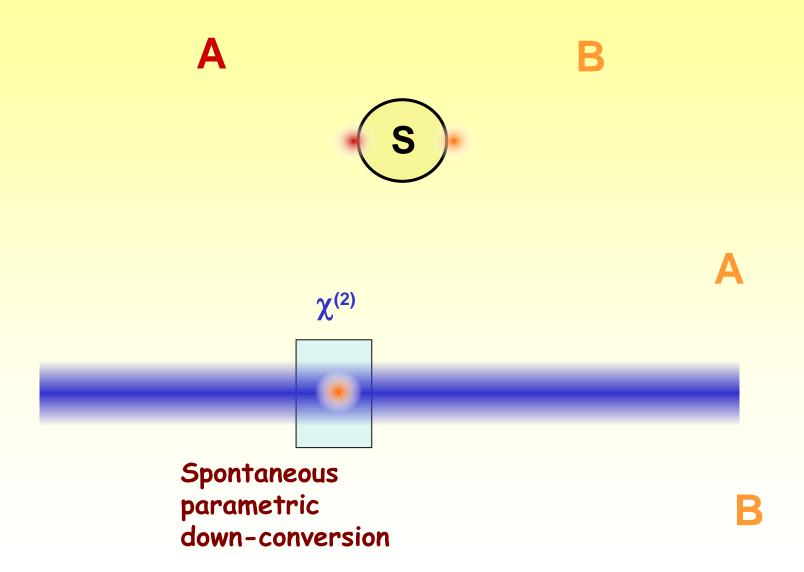




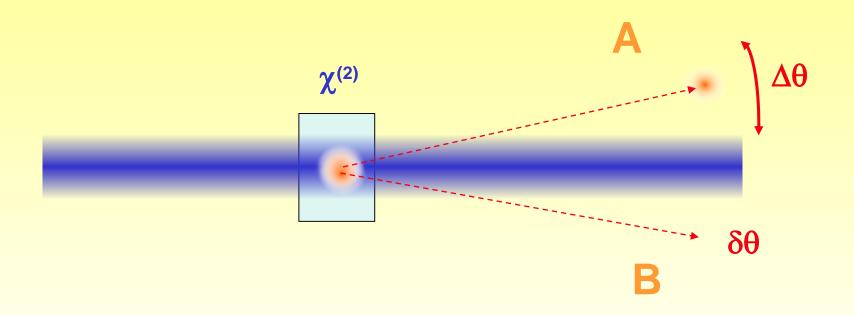




BIPHOTONS



TRANSVERSE ENTANGLEMENT



Entanglement: $\Delta\theta >> \delta\theta$

Uncertainty for one subsystem plus correlations between the two subsystems

TWO-PHOTON AMPLITUDE

$$|\Psi\rangle = \iint d\theta_i d\theta_s F(\theta_i, \theta_s) a^+(\theta_i) a^+(\theta_s) |\operatorname{vac}\rangle$$

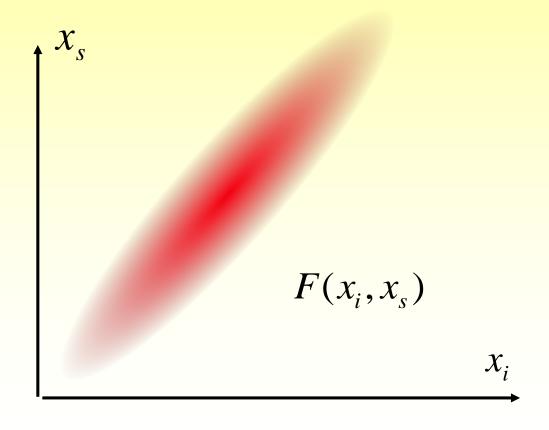
$$heta_s$$
 $F(heta_i, heta_s)$ $heta_i$

$$F(\theta_i, \theta_s) = \sum_{n} \sqrt{\lambda_n} \varphi_n(\theta_i) \chi_n(\theta_s), \sum_{n} \lambda_n = 1, \quad K = \left[\sum_{n} \lambda_n^2\right]^{-1}$$

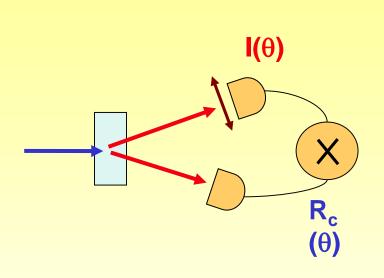
The Schmidt number

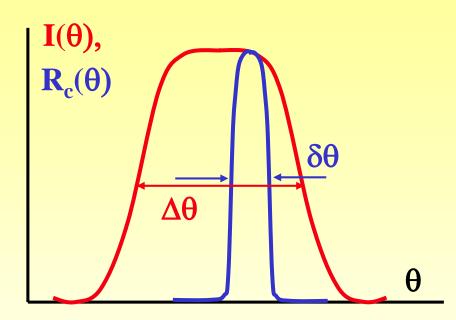
TWO-PHOTON AMPLITUDE: NEAR FIELD

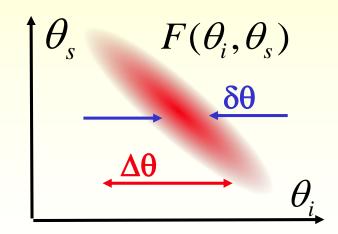
$$|\Psi\rangle = \iint dx_i dx_s F(x_i, x_s) a^+(x_i) a^+(x_s) |\operatorname{vac}\rangle$$



FEDOROV RATIO



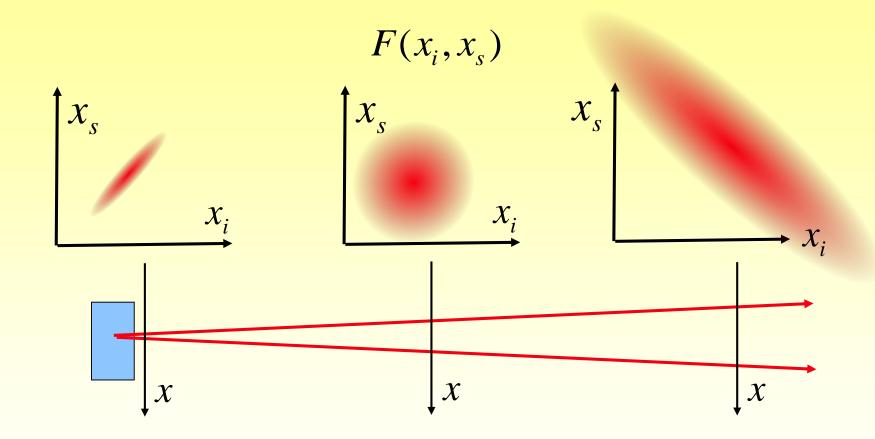




Fedorov ratio: $R = \Delta \theta / \delta \theta$

Valid only for pure states; For Gaussian TPAs coincides with the Schmidt number

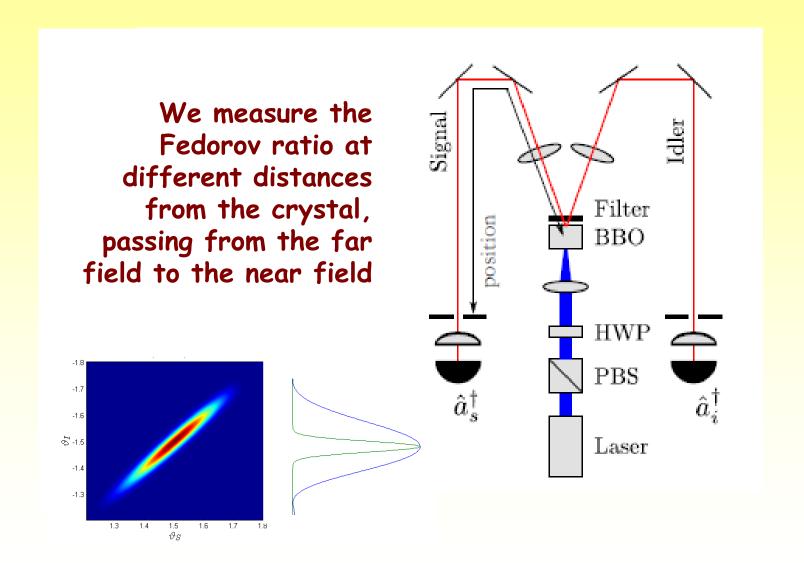
FROM NEAR TO FAR FIELD



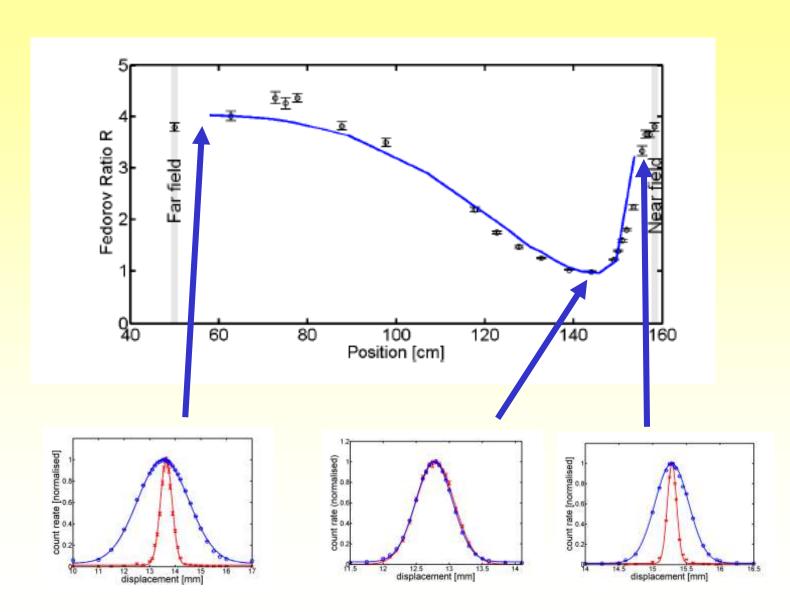
The Fedorov ratio becomes unity between near and far field. But entanglement cannot disappear!

K.W.Chan, J.P. Torres, and J.H.Eberly, PRA 75, 050101 (2007).8

EXPERIMENT

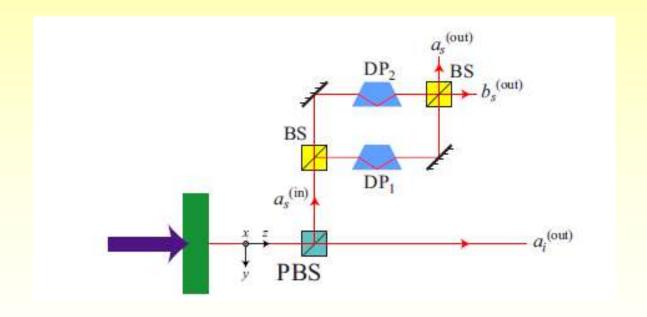


RESULTS: FEDOROV RATIO



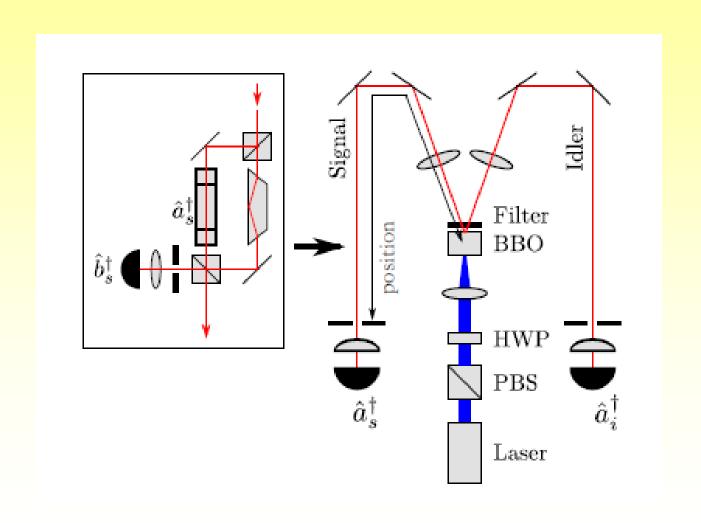
SCHMIDT NUMBER MEASUREMENT

The Schmidt number gives the number of transverse (spatial) modes. The more modes, the less the spatial coherence. Spatial coherence can be measured through interference.



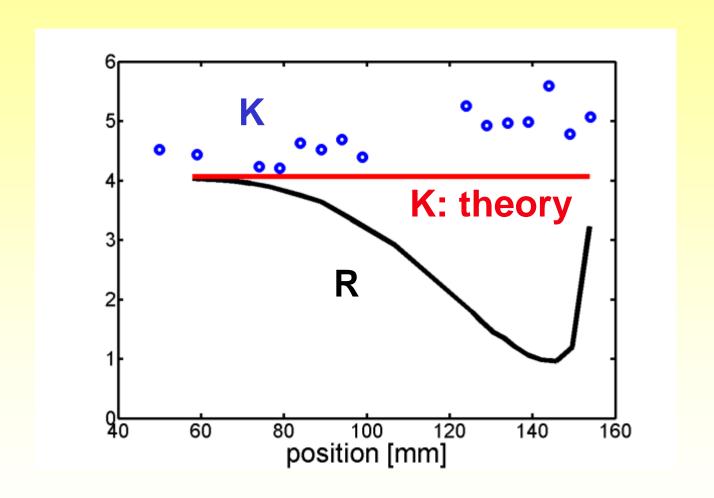
K.W.Chan, J.P. Torres, and J.H. Eberly, PRA 75, 050101 (2007).11

EXPERIMENT



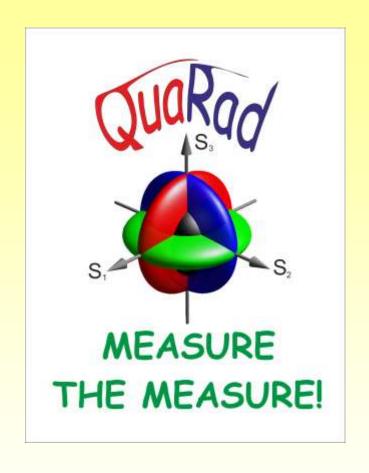
INTERFEROGRAMS

RESULTS: THE SCHMIDT NUMBER



CONCLUSIONS

- 1. The Fedorov ratio can be used as a measure of entanglement only in the near and far field zones.
- 2. As a universal measure of entanglement, the Schmidt number can be used. It can be measured as the inverse visibility of interference. And it does not change as the field propagates from near to far zones.



THANK YOU FOR YOUR ATTENTION!