Nonclassical light in quantum cryptography

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When do we really need nonclassicality of the signal?

How much does it help?

Is it worth the effort?

... is it always good?

Quantum key distribution



The secure key can be distilled if $I_{AB} > I_{BE}$ or $I_{AB} > I_{AE}$.

Lower bound on secure key: $K \ge \max(I_{AB} - I_{BE}, I_{AB} - I_{AE})$





Squeezed states-based protocol:

- Alice generates a Gaussian random variable a
- Alice prepares a squeezed state, displaced by a
- Bob measures a quadrature
- Bases reconciliation
- Error correction, privacy amplification





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Coherent states-based protocol:

- Alice generates two Gaussian random variables {a,b}
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Coherent states-based protocol:

Laser source, quadrature modulation F. Grosshans and P. Grangier. PRL 88, 057902 (2002); F. Grosshans et al., Nature 421, 238 (2003)

Achievements:

80 km [P.Jouguet et al., Nature Photonics 7, 378-381 (2013)]

100 km [D. Huang et al., Sci. Rep. 6, 19201 (2016)]



Two-mode squeezed vacuum state shared between the trusted parties



Two-mode squeezed vacuum: before measurement



Two-mode squeezed vacuum: after homodyne measurement



Two-mode squeezed vacuum: after heterodyne measurement



Allows security analysis based on state purification



CV QKD

Features

- Quadrature encoding & homodyne detection
- Mode description of light
- Gaussian security proofs, optimality of Gaussian attacks
- Covariance matrix formalism (symplectic framework)

$$K = I_{AB} - \chi_{BE} \qquad I_{AB} = \frac{1}{2} \log_2 \frac{V_B}{V_{B|A}} \qquad \chi_{BE} = S(\rho_E) - S(\rho_{E|B})$$

Purification: $S(\rho_E) = S(\rho_{AB})$ $S(\rho_{E|B}) = S(\rho_{A|B})$

Von Neumann entropy:

$$S_{\gamma} = \sum_{i} G\left(\frac{\lambda_{i} - 1}{2}\right) \qquad \qquad G(x) = (x+1)\log_{2}\left(x+1\right) - x\log_{2}x$$

Conditional states:

$$\gamma_E^{x_B} = \gamma_E - \sigma_{BE} (X \gamma_B X)^{MP} \sigma_{BE}^T \qquad \qquad X = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

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Details of security analysis

- E. Diamanti and A. Leverrier, Entropy 17, 6072 (2015) / arXiv:1506.02888
- VCU and R. Filip, Entropy 18, 20 (2016) / arXiv:1601.03105

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Issues

- Gaussian modulation (possibly with a single modulator: VCU, F. Grosshans, Phys. Rev. A 92, 062337 (2015), but still..)
- Channel estimation (can be optimized, L. Ruppert, VCU, and R. Filip, Phys. Rev. A 90, 062310 (2014))

Key distillation



Key distillation: classical algorithms (data manipulation)

- error correction (producing identical data sequences)
- privacy amplification (decoupling Eve form a reference side of the protocol)

Key distillation



Problem: error correction is costly (reduces the mutual information)

$$K = \beta I_{AB} - I_{BE}$$
 where $\beta \in [0,1]$

If we apply modulation independently on the amount of squeezing and optimize it...

1. Squeezed-state protocol tolerates lower post-processing efficiencies



Generalized entanglement-based CV QKD scheme: arbitrary Gaussian modulation of an arbitrarily squeezed state

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1. Squeezed-state protocol tolerates lower post-processing efficiencies



Maximum tolerable versus signal squeezing noise upon limited post-processing efficiency (from top to bottom: $\beta = 0.8, 0.6, 0.4, 0.2$) [VCU and R. Filip, New J. Phys. 13, 113007 (2011)]

If we apply modulation independently on the amount of squeezing and optimize it...

2. Squeezed-state protocol can tolerate more noise/loss than any coherent-state CV QKD protocol



Sketch of the experiment, performed at DTU in Lyngby (group of Ulrik Andersen)

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Maximum tolerable channel noise versus modulation (left) and maximum tolerable noise for given channel noise (right) for optimized coherent-state protocol (grey area) and squeezed-state (theory dashed lines + experimental points)

[L. Madsen, VCU, M. Lassen, R. Filip, U. Andersen, Nat. Comm. 3, 1083 (2012)]

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Correlation of the outputs of a beamsplitter: $C_{BE} \propto V_S - V_E = (V + \sigma) - 1 = 0$ Condition for vanishing of Holevo bound: $\sigma = 1 - V$

$$\begin{pmatrix} V_{s} & 0 & 0 & 0 \\ 0 & V_{s}^{(p)} & 0 & 0 \\ 0 & 0 & V_{E} & 0 \\ 0 & 0 & 0 & V_{E}^{(p)} \end{pmatrix} \rightarrow \begin{pmatrix} \eta V_{s} + (1 - \eta) V_{E} & 0 & \sqrt{\eta (1 - \eta)} (V_{E} - V_{s}) & 0 \\ 0 & \eta V_{s}^{(p)} + (1 - \eta) V_{E}^{(p)} & 0 & \sqrt{\eta (1 - \eta)} (V_{E}^{(p)} - V_{s}^{(p)}) \\ \sqrt{\eta (1 - \eta)} (V_{E} - V_{s}) & 0 & \eta V_{E} + (1 - \eta) V_{s} & 0 \\ 0 & \sqrt{\eta (1 - \eta)} (V_{E}^{(p)} - V_{s}^{(p)}) & 0 & \eta V_{E}^{(p)} + (1 - \eta) V_{s}^{(p)} \end{pmatrix}$$

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...but if a pre-modulation lossy side channel on the sender side is present...

Squeezed-state protocol is more sensitive to the side-channel loss



Lossy side channel prior to state modulation

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Squeezed-state protocol is more sensitive to the side-channel loss



Key rate vs distance without (solid lines) and with a 50% pre-modulation side channel. Blue: squeezed states (0.1, 0.5 SNU), orange: coherent-state protocol (no effect). [I. Derkach, VCU, and R. Filip, PRA 93, 032309 (2016) + in preparation]



Typical noise model used in CV QKD and parametrized by a mean photon number



The same noise model applied to DV QKD protocol



DV security analysis

$$K^{(BB84)} = p_{exp} \max[0, 1 - 2H(Q)]$$

$$K^{(6state)} = p_{exp} \max\left[0, 1 - F(Q)\right]$$

$$F(Q) = -\left(1 - \frac{3Q}{2}\right)\log_2\left(1 - \frac{3Q}{2}\right) - \frac{3Q}{2}\log_2\frac{Q}{2}$$

[B. Kraus, N. Gisin, and R. Renner, Phys. Rev. Lett. 95, 080501 (2005); R. Renner, N. Gisin, and B. Kraus, Phys. Rev. A 72, 012332 (2005)]

Evaluation of QBER

Click on a "right" detector along with *k* noise photons on the "right" and *I* noise photons on the "wrong" detector:

$$p_+(k,l) = T\pi_k(T)\pi_l(T)$$

Click on a "wrong" detector: $p_{-}(k,l) = (1-T)\pi_{k}(T)\pi_{l}(T)$

Where
$$\pi_k(T) = \sum_{n=k}^{\infty} p_n(\mu) \binom{n}{k} (1-T)^k T^{n-k}$$

[VCU, M. G. A. Paris, Phys. Lett. A 374, 1342 (2010)]

Expected probability of accepting a given event:

$$p_{exp} = \sum_{k=0}^{\infty} p_+(k,0) + \sum_{k=1}^{\infty} p_-(k,0) + \sum_{l=1}^{\infty} p_-(0,l)$$
 then $Q = \frac{\sum_{l=1}^{\infty} p_-(0,l)}{p_{exp}}$



Comparison between robustness to noise in DV and CV

Analytical result for CV:

$$\mu_{\max}(T) = \exp[1 + W_{-1}(-T/e)]$$

Analytical result for DV:

$$\mu_{\max}^{DV}(T) = \frac{TQ_{\text{th}}}{1 - 2Q_{\text{th}}}$$
$$(Q_{th} \approx 12.6\% \text{ for BB84})$$



Requirements on nonclassicality of the sources for CV and DV in noisy channels



How good shall be the single-photon DV source to beat any CV protocol

[M. Lasota, R. Filip, VCU, arXiv:1602.03122]

...and much more:

- For CV squeezing is also helpful in fluctuating channels (test in progress)
- Squeezed states improve channel estimation [L. Ruppert, VCU, and R. Filip, Phys. Rev. A 90, 062310 (2014)]
- For DV in noisy channels non-Gaussianity can indicate the suitability of a channel to QKD (nonclassicality is not sufficient) [arXiv:1603.06620]

...and much more:

- For CV squeezing is also helpful in fluctuating channels (test in progress)
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What next?

- Incorporate decoy-state DV QKD
- Entanglement-based schemes
- Better witnesses for channel verification

To sum up:

- In CV QKD nonclassicality and optimal use of resources is helpful, unless side-channel loss on the sender side is present (and channel noise is low).
- Single-photon DV QKD can be more robust to channel noise than any squeezed-state CV QKD.

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Thank you for attention!

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